
Erratum

Stochastic Model of a One-Dimensional Fluid¹

Sezar Fesciyan² and H. L. Frisch²

Received April 4, 1975

The derivation of z given on p. 157 of this paper¹ was based on the standard thermodynamic relation $\partial \ln z / \partial p = 1/\alpha$. For explicitly density-dependent potentials this relation is not quite true. From the expression for the Gibbs free energy

$$G(p, N) = -N \ln \int_0^\infty e^{-px} e^{-\phi(x,p)} dx$$

we see that

$$dG(p, N) = \left[L + N \left\langle \frac{\partial \phi}{\partial p} \right\rangle \right] dp + \left[\mu + \int^p \left\langle \frac{\partial \phi}{\partial p} \right\rangle dp \right] dN$$

where

$$\left\langle \frac{\partial \phi}{\partial p} \right\rangle = \int_0^\infty \frac{\partial \phi(x, p)}{\partial p} \alpha P_0^\alpha(x) dx$$

The total chemical potential has a contribution coming from "polarization" effects. The activity z derived in this paper really represents e^μ with μ defined

¹ This paper appeared in *J. Stat. Phys.* **12**(2):153 (1975).

² Department of Chemistry, State University of New York at Albany, Albany, New York.

as above. This also affects the proof of density dependence on p. 161. The expression following Eq. (20) should instead read

$$\left[L \left(\frac{\partial p}{\partial \alpha} \right) + N \frac{(1 - \partial p / \partial \alpha)}{\alpha} - N \left\langle \frac{\partial \phi}{\partial \alpha} \right\rangle \right] P(N, L) - \frac{N + 1}{\alpha} P(N + 1, L)$$

which leaves the required proof inconclusive. However, the fact that the pair potential is explicitly density dependent can be proved easily by observing that in the region $\sigma < x < 2\sigma$ we have $P_0^{(\infty)}(x) = 1$, yielding $\partial \phi / \partial x = -p$, which is clearly density dependent.